

```

rstN[a_, s_, t_, N_] := E^(I ((1 - a) s - (2 - a) t) / ((1 - a) N));
ystN[a_, s_, t_, N_] := E^(I (t - s) / ((1 - a) N));
zstN[a_, s_, t_, N_] := E^(I * t / N);
costthetaa[a_, r_, y_, z_] := betaa[a, r, y, z] / (2 * a * z * tau[a, a, r, y, z]);
Clear[thetaa];
thetaa[a_, s_, t_, N_] := I * Sqrt[(1 - a) s^2 + a * t^2] / N;

```

\*\*\*\* VERIFY first expression in (1.92) for  $\cos \theta[r_{s,t,N}, y_{s,t,N}, z_{s,t,N}]$

```

Collect[Normal[
  Series[Normal[Series[costthetaa[a, rstN[a, s, t, N], ystN[a, s, t, N], zstN[a, s, t, N]] - 
    (1 + (1/2) ((1 - a) s^2 + a * t^2) / N^2), {t, 0, 2}], {s, 0, 2}]], N]

```

$$\frac{\frac{3 i a s^2 t}{2 (-1+a)} - \frac{2 i a^2 s^2 t}{-1+a} + \frac{i a^3 s^2 t}{2 (-1+a)} + \frac{1}{2} i a^2 s t^2}{N^3} - \frac{a s^2 t^2}{-1+a} + \frac{17 a^2 s^2 t^2}{4 (-1+a)} - \frac{5 a^3 s^2 t^2}{2 (-1+a)}$$

\*\*\*\* VERIFY second expression in (1.92) for  $\theta[r_{s,t,N}, y_{s,t,N}, z_{s,t,N}]$

```

Collect[Normal[
  Series[Normal[Series[Cos[thetaa[a, s, t, N]] - (1 + (1/2) ((1 - a) s^2 + a * t^2) / N^2), 
    {t, 0, 3}], {s, 0, 3}]], N]

```

$$\frac{s^2 (a t^2 - a^2 t^2)}{12 N^4}$$

\*\*\* CONCLUSION: thetaa[a, r, s, t, N] yields the correct expansion of  $\cos \theta[r_{s,t,N}, y_{s,t,N}, z_{s,t,N}]$  through order  $N^{-3}$ , so there is no non-zero term of order  $N^{-2}$  in the expansion of  $\theta[r_{s,t,N}, y_{s,t,N}, z_{s,t,N}]$ .

Since the model is homogeneous, we need only apply the first line of (1.78) with  $f := N$  to obtain

$$w_N^*(a) = (\sqrt{x_a} \sin \theta)^{-1} [az\tau_a]^N \{\sin N\theta - \sqrt{x_a} w_0^*(a) \sin(N-1)\theta\}. \quad (1.93)$$

Expand  $\sin(N-1)\theta = \mathbf{s} \cos \theta - \mathbf{c} \sin \theta$ , for  $\mathbf{s} := \sin N\theta$  and  $\mathbf{c} := \cos N\theta$ . Put  $\Lambda := (\sin \theta)^{-1} [az\tau_a]^{N-1}$ . After direct calculation we find  $1 - \sqrt{x_a} w_0^*(a) = 1 - a + O(N^{-1})$ . Therefore by (1.93) we have

$$\frac{w_N^*(a)}{\Lambda} = \mathbf{s} - \sqrt{x_a} w_0^*(a)(\mathbf{s} \cos \theta - \mathbf{c} \sin \theta) = (1-a)\mathbf{s} + \frac{O(1)}{N^1}. \quad (1.94)$$

Note that there is no cancellation of the order 1 term in (1.94). Now plug (1.94) into (1.37) to obtain

$$g_{0,N} = \frac{\omega_a}{a(2-a)} r z \tau_a^{-1} \frac{(N - (N-1)a) \sin \theta}{(1-a)\mathbf{s} + O(N^{-1})}.$$

Finally apply the asymptotic expression for  $\theta$  in (1.92) and let  $N \rightarrow \infty$ .  $\square$

\*\*\*\*\* EXPAND \$ (1- \sqrt{x\_b} w\_0^\*(b)) [r\_{s,t,N}, y\_{s,t,N}, z\_{s,t,N}] \$ to pass from (1.93) to (1.94).

```
Collect[Normal[
  Series[Normal[Series[1 - a * zN[a, s, t, N] * tau[a, a, rstN[a, s, t, N], ystN[a, s, t, N],
    zN[a, s, t, N]] * w0stara[a, rstN[a, s, t, N],
    ystN[a, s, t, N], zN[a, s, t, N]], {t, 0, 1}]], {s, 0, 1}]], N]
```

$$1 - a + \frac{i a s - i a^2 s - 2 i a t + i a^2 t}{N} + \frac{3 a^2 s t - 2 a^3 s t}{N^2}$$

*Proof (Theorem 1.1.6).* By the same reasoning given at the outset of the proof of Theorem 1.1.1, we may assume that  $\mathbf{X}_0 = 0$ . By the fact that the absolute value process starts afresh at the end of each excursion, we have that  $1 + \mathcal{M}_N$  is a standard geometric random variable with success probability  $P(\mathbf{H} \geq N)$ . Thus

$$P(\mathcal{M}_N = \nu) = [P(\mathbf{H} < N)]^\nu P(\mathbf{H} \geq N), \quad \nu = 0, 1, 2, \dots \quad (1.95)$$

Let  $\mathbf{L}_N$ ,  $\mathbf{R}_N$ , and  $\mathbf{V}_N$ , respectively, be random variables for the number of steps, runs, and short runs, in an excursion, given that the height of the excursion is at most  $N - 1$ . Therefore, in distribution, we may write:

$$\mathcal{R}_N = \sum_{\nu=0}^{\mathcal{M}_N} \mathbf{R}^{(\nu)}, \quad \mathcal{V}_N = \sum_{\nu=0}^{\mathcal{M}_N} \mathbf{V}^{(\nu)}, \quad \mathcal{L}_N = \sum_{\nu=0}^{\mathcal{M}_N} \mathbf{L}^{(\nu)}, \quad (1.96)$$

where  $\mathbf{R}^{(1)}, \mathbf{R}^{(2)}, \dots$ ;  $\mathbf{V}^{(1)}, \mathbf{V}^{(2)}, \dots$ ; and  $\mathbf{L}^{(1)}, \mathbf{L}^{(2)}, \dots$ , respectively, are sequences of independent copies of  $\mathbf{R}_N$ ,  $\mathbf{V}_N$ , and  $\mathbf{L}_N$ . Since the random variables  $\mathbf{R}_N$ ,  $\mathbf{V}_N$ , and  $\mathbf{L}_N$  already have built into their definitions the condition  $\{\mathbf{H} \leq N - 1\}$ , the probability generating function  $K_{N-1} = E\{r^{\mathbf{R}_N} y^{\mathbf{V}_N} z^{\mathbf{L}_N}\}$

is calculated by Theorem 1.3.16. Thus by (1.95), and by calculating a geometric sum there holds:

$$E\{r^{\mathcal{R}_N}y^{\mathcal{V}_N}z^{\mathcal{L}_N}u^{\mathcal{M}_N}\} = \sum_{\nu=0}^{\infty} P(\mathcal{M}_N = \nu) (uK_{N-1})^\nu = \frac{P(\mathbf{H} \geq N)}{1-uP(\mathbf{H} < N)K_{N-1}[r,y,z]}. \quad (1.97)$$

We define  $(r_N, y_N, z_N)$  by (1.73), and also set  $u_N := e^{-ita(b-a)/[(1-a)(1-b)N]}$ . By (1.9), it suffices to show that  $\lim_{N \rightarrow \infty} E\{e^{it(1+1/N)\mathcal{X}_{N+1}}\} = \hat{\psi}(t)/\hat{\varphi}(t)$ ; see (1.102). We define  $\theta_1$  and  $\theta_2$  by (1.74), so that also (1.75)–(1.76) hold. By the statement of Theorem 1.3.16 we must replace the calculation of  $\bar{w}_{0,N}$ , starting with (1.77), with instead  $\bar{w}_{1,N+1}$ . However, by (1.53), (1.62), and (1.77), the difference in the two calculations is simply accounted for by replacing  $f$  by  $f-1$  in the calculation of  $\bar{w}_{0,N}$ , because  $j$  in (1.62) for  $\bar{w}_{1,N+1}$  is determined by  $j = N+1-f = N-(f-1)$ , so  $\frac{1}{\Lambda_2}\bar{w}_{1,N+1} = d_1(f')q^* + d_2(f')w^*$  with  $f' := f-1$  in place of  $f$  in both (1.53) and (1.83). This is reflected by the fact that, by Lemma 1.3.8,  $\bar{w}_{1,N+1} = \bar{w}_{N,0}$ . We must now also calculate  $\bar{q}_N = d_{q,1}(f)q_{N-f+1}^*(b) + d_{q,2}(f)w_{N-f+1}^*(b)$  given by Lemma 1.3.15, with  $d_{q,j}(f) = \mu_{j,1}q_{f-1}^*(a) + \mu_{j,2}q_f^*(a)$ ,  $j = 1, 2$ , defined by (1.61) in the proof of Lemma 1.3.15. In summary,  $f' = f-1$  yields  $(\dagger) \bar{q}_N = d_{q,1}(f'+1)q_{N-f'}^*(b) + d_{q,2}(f'+1)w_{N-f'}^*(b)$ . Thus, because we simply replace  $f$  by  $f-1$  in the required substitutions, and since  $f \sim \eta N$ , we will not change the name of  $f$ . With this understanding, we may use the calculation of  $\bar{w}_{0,N}$  in (1.77)–(1.89) verbatim in place of the calculation of  $\bar{w}_{1,N+1}$ , and we will do this without changing the names of  $e_j$ ,  $q^*$ ,  $w^*$  and  $\Lambda_j$ ; see (1.79) and (1.83). Further with this understanding, by  $(\dagger)$ , with  $f$  now recouping the role of  $f'$ , and with  $q^*$  and  $w^*$  defined by (1.83), we have  $\frac{1}{\Lambda_2}\bar{q}_N = d_{q,1}q^* + d_{q,2}w^*$  for

$$d_{q,j} := \mu_{j,1}q_f^*(a) + \mu_{j,2}q_{f+1}^*(a). \quad (1.98)$$

Here by (1.13), (1.41) and (1.74), in analogy with (1.78), we have

$$\begin{aligned} q_f^*(a) &= 2i\alpha_a^{-1}[az\tau_a]^f \{y^2 \sin f\theta_1 - \sqrt{x_a}q_0^*(a) \sin(f-1)\theta_1\} \\ q_{f+1}^*(a) &= 2i\alpha_a^{-1}[az\tau_a]^f \sqrt{x_a} \{y^2 \sin(f+1)\theta_1 - \sqrt{x_a}q_0^*(a) \sin f\theta_1\}. \end{aligned} \quad (1.99)$$

Denote  $e_{q,j} := d_{q,j}/\Lambda_1$ , and  $f_+ = f+1$  and  $f_-$  as before. Therefore, by (1.98)–(1.99),

$$e_{q,j} = (y^2\mu_{j,1} - \mu_{j,2}xq_0^*) \sin f\theta_1 + \sqrt{x} \{y^2\mu_{j,2} \sin f_+\theta_1 - \mu_{j,1}q_0^* \sin f_-\theta_1\}, \quad (1.100)$$

where  $x = x_a$  and  $q_0^* = q_0^*(a)$ . Rewrite (1.100) by applying the notations (1.81). Thus  $e_{q,j}$  is written, with dependence on  $a$  suppressed, by

$$(y^2\mu_{j,1} - \mu_{j,2}xq_0^*)\mathbf{s}_1 + \sqrt{x} \{y^2\mu_{j,2}(\mathbf{s}_1 \cos \theta_1 + \mathbf{c}_1 \sin \theta_1) - \mu_{j,1}q_0^*(\mathbf{s}_1 \cos \theta_1 - \mathbf{c}_1 \sin \theta_1)\} \quad (1.101)$$

In summary, by (1.98), we have  $\bar{q}_N/(\Lambda_1\Lambda_2) = e_{q,1}q^* + e_{q,2}w^*$ , for  $e_{q,j}$  in (1.101), and  $\Lambda_j$  defined by (1.79) and (1.83).

To guide the asymptotic expansions of (1.101) we rewrite (1.97) by substituting the last line of the proof of Theorem 1.3.16:

$$E\{e^{it(1+1/N)\mathcal{X}_{N+1}}\} = \frac{P(\mathbf{H} \geq N+1)\bar{w}_{1,N+1}}{\bar{w}_{1,N+1} - (1-a)u_N r_N^2 z_N^2 \bar{q}_N}. \quad (1.102)$$

It turns out that there is a cancellation in the order of the denominator of (1.102). That is, the leading order of each of  $\bar{w}_{1,N+1}/(\Lambda_1 \Lambda_2)$  and  $\bar{q}_N/(\Lambda_1 \Lambda_2)$  will be some order 1 trigonometric factor times  $it/N$ ; in fact there holds  $(1-a)\bar{q}_N/\bar{w}_{1,N+1} \sim 1$ , as  $N \rightarrow \infty$ . Define

$$\Delta_N := \bar{w}_{1,N+1} - (1-a)u_N r_N^2 z_N^2 \bar{q}_N. \quad (1.103)$$

By direct calculation we will establish that  $\Delta_N/(\Lambda_1 \Lambda_2) = O(N^{-2})$ , and we find the exact coefficient of the order  $N^{-2}$  term.

For the asymptotics of (1.101) we may still treat  $\sin \theta_1 = \theta_1 + O(N^{-3})$  by (1.76), but must render precisely the  $O(N^{-2})$  term in  $\cos \theta_1 = 1 + O(N^{-2})$  of (1.75). In an appendix to [15], we display the many terms of the book-keeping method for this problem. For the present, we simply exhibit the asymptotics of (1.103) obtained by machine computation with  $\sin \theta_j$  substituted by the corresponding order  $1/N$  term of (1.76):

$$\frac{\Delta_N}{\Lambda_1 \Lambda_2} = \frac{1}{(1-a)(1-b)} \frac{t^2}{N^2} \{-ab\sigma_1\sigma_2\mathbf{c}_1\mathbf{c}_2 - a\sigma_1(a-b)^2\mathbf{c}_1\mathbf{s}_2 + a^2\sigma_1^2\mathbf{s}_1\mathbf{s}_2\} + \frac{O(1)}{N^3}. \quad (1.104)$$

Finally we compute the limit of the ratio (1.102) by the asymptotic relations (1.74), and by (1.89) and (1.104). Thus, because by (1.64) and Proposition 1.3.2 we have that  $P(\mathbf{H} \geq N+1) \sim C_{a,b}N^{-1}$  for  $C_{a,b} = ab/[(1-\eta)a + \eta b - ab]$ , we find  $E\{e^{it(1+1/N)\mathcal{X}_{N+1}}\}$  is asymptotic to

$$\begin{aligned} & C_{a,b}N^{-1} \left\{ [a\sigma_1\mathbf{c}_1\mathbf{s}_2 + b\sigma_2\mathbf{s}_1\mathbf{c}_2 + (b-a)^2\mathbf{s}_1\mathbf{s}_2] \frac{it}{N} + O(N^{-2}) \right\} \\ & \div \left\{ \frac{1}{(1-a)(1-b)} [-ab\sigma_1\sigma_2\mathbf{c}_1\mathbf{c}_2 - a\sigma_1(a-b)^2\mathbf{c}_1\mathbf{s}_2 + a^2\sigma_1^2\mathbf{s}_1\mathbf{s}_2] \frac{t^2}{N^2} + \frac{O(1)}{N^3} \right\}. \end{aligned} \quad (1.105)$$

But, as in the proof of Theorem 1.1.1 we have  $\mathbf{c}_j \sim \cosh(\kappa_j t)$ , and  $\mathbf{s}_j \sim i \sinh(\kappa_j t)$ ,  $j = 1, 2$ . Therefore, with  $\tilde{C}_{a,b} := (1-a)(1-b)C_{a,b}$ , we obtain that (1.105) has the following limit as  $N \rightarrow \infty$ , where we refer to (1.3) and statement of Theorem 1.1.6 for the definitions of  $\hat{\varphi}(t)$  and  $\hat{\psi}(t)$ :

$$\lim_{N \rightarrow \infty} E\{e^{it(1+1/N)\mathcal{X}_N}\} = \frac{\tilde{C}_{a,b}}{t} \times \frac{(b\kappa_1\sigma_2 + a\kappa_2\sigma_1)t}{\hat{\varphi}(t)} \times \frac{\hat{\psi}(t)}{ab\sigma_1\sigma_2}.$$

We have  $\tilde{C}_{a,b} = ab\sigma_1\sigma_2/(a\sigma_1\kappa_2 + b\sigma_2\kappa_1)$ , so the proof is complete.  $\square$

\*\*\*\*\* DEFINE e\_{q,1}, e\_{q,2}, with auxiliary variables s1 and c1sin as in the proof of Theorem 1.1.1.

Also define the centering term  $u_N$ .

```

eq1[a_, b_, r_, y_, z_, s1_, c1sin_] :=
  (y^2 * m11[a, b, r, y, z] - m12[a, b, r, y, z] * x[a, r, y, z] * q0stara[a, r, y, z]) s1 +
  a * z * tau[a, a, r, y, z] * (y^2 * m12[a, b, r, y, z] (s1 * costheta1[a, r, y, z] + c1sin) -
  m11[a, b, r, y, z] * q0stara[a, r, y, z] (s1 * costheta1[a, r, y, z] - c1sin));
eq2[a_, b_, r_, y_, z_, s1_, c1sin_] :=
  (y^2 * m21[a, b, r, y, z] - m22[a, b, r, y, z] * x[a, r, y, z] * q0stara[a, r, y, z]) s1 +
  a * z * tau[a, a, r, y, z] * (y^2 * m22[a, b, r, y, z] (s1 * costheta1[a, r, y, z] + c1sin) -
  m21[a, b, r, y, z] * q0stara[a, r, y, z] (s1 * costheta1[a, r, y, z] - c1sin));
uN[a_, b_, t_, N_] := E^((-I * t * a (b - a)) / ((1 - a) (1 - b) N));

```

\*\*\*\*\* AUTOMATIC CALCULATION of the asymptotics of  $\Delta_N / (\Lambda_1 \Lambda_2)$  [  $r_N, y_N, z_N, u_N$  ] for (1.104)

```

Collect[
  Factor[Normal[Series[e1[a, b, rN[a, b, t, N], yN[a, b, t, N], zN[a, b, t, N], s1, c1 *
    sintheta1asymp[a, b, t, sig1, N]] * (s2 * s2qstar[a, b, rN[a, b, t, N],
    yN[a, b, t, N], zN[a, b, t, N]] + c2 * sintheta2asymp[a, b, t, sig2, N] *
    c2sinqstar[a, b, rN[a, b, t, N], yN[a, b, t, N], zN[a, b, t, N]]) +
    e2[a, b, rN[a, b, t, N], yN[a, b, t, N], zN[a, b, t, N], s1,
    c1 * sintheta1asymp[a, b, t, sig1, N]] *
    (s2 * s2wstar[a, b, rN[a, b, t, N], yN[a, b, t, N], zN[a, b, t, N]] +
    c2 * sintheta2asymp[a, b, t, sig2, N] *
    c2sinwstar[a, b, rN[a, b, t, N], yN[a, b, t, N], zN[a, b, t, N]]) -
    (1 - a) uN[a, b, t, N] * rN[a, b, t, N]^2 * zN[a, b, t, N]^2 *
    (eq1[a, b, rN[a, b, t, N], yN[a, b, t, N], zN[a, b, t, N], s1,
      c1 * sintheta1asymp[a, b, t, sig1, N]] * (s2 * s2qstar[a, b, rN[a, b, t, N],
        yN[a, b, t, N], zN[a, b, t, N]] + c2 * sintheta2asymp[a, b, t, sig2, N] *
        c2sinqstar[a, b, rN[a, b, t, N], yN[a, b, t, N], zN[a, b, t, N]]) +
      eq2[a, b, rN[a, b, t, N], yN[a, b, t, N], zN[a, b, t, N], s1,
        c1 * sintheta1asymp[a, b, t, sig1, N]] * (s2 * s2wstar[a, b, rN[a, b, t, N],
          yN[a, b, t, N], zN[a, b, t, N]] + c2 * sintheta2asymp[a, b, t, sig2, N] *
          c2sinwstar[a, b, rN[a, b, t, N], yN[a, b, t, N], zN[a, b, t, N]]))
    , {t, 0, 2}]], N]

```

$$-\frac{1}{(-1+a)(-1+b)N^2} a \left( -a^2 s1 s2 + 2 a^2 b s1 s2 - a b^2 s1 s2 + a^2 c1 s2 \text{sig1} - 2 a b c1 s2 \text{sig1} + b^2 c1 s2 \text{sig1} + b c1 c2 \text{sig1} \text{sig2} \right) t^2$$

```
Factor[-a (-a^2 s1 s2 + 2 a^2 b s1 s2 - a b^2 s1 s2) / (sigma1[a, b]^2)]
```

```
a^2 s1 s2
```

```
Factor[-a (a^2 c1 s2 sig1 - 2 a b c1 s2 sig1 + b^2 c1 s2 sig1)]
```

```
- a (a - b)^2 c1 s2 sig1
```

```
Factor[-a (b c1 c2 sig1 sig2)]
```

```
- a b c1 c2 sig1 sig2
```

$$\begin{aligned} \text{*** CONCLUSION: we already distributed the factor } (-a) \text{ from } & -\frac{1}{(-1+a)(-1+b)N^2} a. \text{ Hence } \Delta_N / (\Lambda_1 \Lambda_2) [rN, yN, zN, uN] \\ &= [a^2 \operatorname{sig1}^2 s1 s2 - a (a - b)^2 c1 s2 \operatorname{sig1} - a b c1 c2 \operatorname{sig1} \operatorname{sig2}] \frac{t^2}{(-1+a)(-1+b)N^2} + O(N^{-3}). \end{aligned}$$

**Corollary 1.4.1.** Assume  $a = b$ . Define

$$Z_1 = \frac{1}{N} \left( \mathcal{R}_N - \frac{1}{(1-a)} \mathcal{V}_N + a \mathcal{M}_N \right); Z_2 = \frac{1}{N} \left( \mathcal{L}_N - \frac{1}{(1-a)} \mathcal{R}_N + \frac{a}{1-a} \mathcal{M}_N \right) - Z_1.$$

$$\text{Then, } \lim_{N \rightarrow \infty} E\{e^{i(sZ_1+tZ_2)}\} = \frac{\tanh(\sqrt{1-a}s^2+at^2)}{\sqrt{1-a}s^2+at^2}.$$

*Proof (Corollary 1.4.1).* One simplifies the lines of proof Theorem 1.1.6. We leave details in an appendix to [15].  $\square$

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## 1.5 Appendix

We display the explicit asymptotic expansions through order  $N^{-2}$  following from (1.82).

$$[\mathbf{s}_1](e_1) = -(1-a)(1-b) + 2(1-ab)\frac{it}{N} + \frac{2-a^2+a^3-3ab-a^2b+2a^2b^2}{(1-a)(1-b)}\frac{t^2}{N^2} + \frac{O(1)}{N^3}. \quad (1.106)$$

$$[\mathbf{s}_1](e_2) = 1 - a - \frac{a(b-a)}{1-b}\frac{it}{N} - \frac{a(-a-5a^2+2a^3+8ab+4a^2b+2b^2-10ab^2)}{2(1-a)(1-b)^2}\frac{t^2}{N^2} + \frac{O(1)}{N^3}. \quad (1.107)$$

For the proof of Theorem 1.1.1 we do not require the explicit order  $N^{-2}$  terms in (1.106)–(1.107), but we display them because they are implicitly required for the proof of Theorem 1.1.6. Again by (1.82) and direct calculation,

$$[\mathbf{c}_1 \sin \theta_1](e_1) = \frac{O(1)}{N^2}; \quad [\mathbf{c}_1 \sin \theta_1](e_2) = a + \frac{a(b-a)}{1-b}\frac{it}{N} + \frac{O(1)}{N^2}. \quad (1.108)$$